

*Short note*

## New states in $^{44,46}\text{Ar}$ isotopes from deep-inelastic heavy ion reaction studies

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**Abstract.** In-beam  $\gamma$ -rays in neutron-rich *sdfp* shell nuclei, produced in deep-inelastic collisions of  $^{48}\text{Ca}$  projectiles on  $^{48}\text{Ca}$  target nuclei, have been studied using Ge multidetector array EU-ROBALL III. New yrast states in heavy argon isotopes have been identified.

**PACS.** 21.60.Cs Shell model – 23.20.Lv Gamma transitions and level energies – 27.40.+z  $39 \leq A \leq 58$  – 25.70.Lm Strongly damped collisions

Neutron-rich nuclei close to the  $N=28$  shell gap have recently attracted a particular interest triggered by a possible existence of anomalies in the shell closure. Indeed, the measured half-lives of  $\beta$ -decaying  $^{44}\text{S}$  and  $^{45-47}\text{Cl}$  nuclei were found much shorter than those predicted by calculations, which might indicate an unexpected shape transition in that region. On the basis of self-consistent mean-field calculations Werner et al. [1] suggested, that the major  $N=28$  shell gap disappears when approaching  $Z=16$  and the energy surfaces become very soft with close lying and shallow minima corresponding to different deformations.

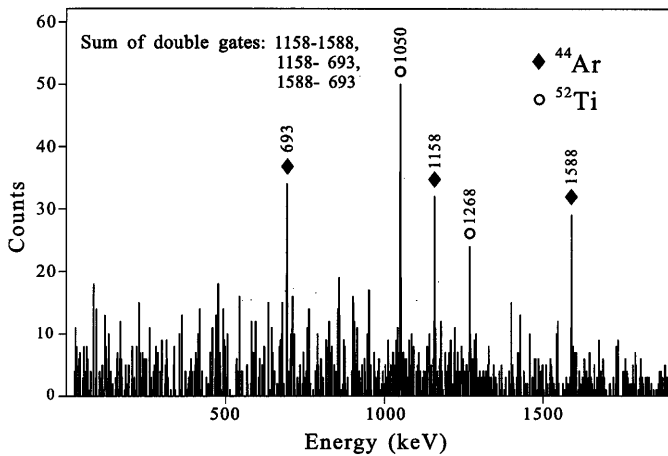
Some of these exotic nuclei have been investigated using a technique of intermediate-energy Coulomb excitation of radioactive beams. Energies of the  $2_1^+$  states and  $B(E2; 0_{g.s.}^+ \rightarrow 2_1^+)$  values in  $^{40,42}\text{S}$  and in  $^{44,46}\text{Ar}$  [2] have been measured. Similar study has been performed with the relativistic radioactive ion beam taking advantage of larger cross-sections for Coulomb excitations at those energies [3]. In  $^{44}\text{Ar}$  candidates for higher lying excitations: a second  $2^+$  level and a  $4^+$  state, were suggested.

In a series of recent experiments we have shown that the yrast spectroscopy of hard-to-reach neutron-rich nuclei, populated in heavy-ion multinucleon transfer reactions (at bombarding energies of  $\sim 15\%$  above Coulomb barrier), can be studied very successfully in  $\gamma$ - $\gamma$  thick target coincidence measurements (e.g. [4,5]). In such deep-inelastic collisions the population of yrast states at moderate spins is strongly favoured.

A possibility to access new nuclei in the  $N=20-28$  neutron-rich region is offered by the  $^{48}\text{Ca}$  beam. We performed an experiment at the Tandem accelerator in Laboratori Nazionali di Legnaro bombarding a target of  $0.74 \text{ mg/cm}^2$   $^{48}\text{Ca}$  (backed by  $40 \text{ mg/cm}^2$  of evaporated  $^{208}\text{Pb}$ ) with a beam of  $140 \text{ MeV}$   $^{48}\text{Ca}$  ions. The  $\gamma$ - $\gamma$  coincidences were collected with the Euroball III array. Fusion-evaporation was a main reaction channel whereas multinucleon transfer processes, leading to nuclei from the vicinity of  $^{48}\text{Ca}$ , accounted for less than  $1\%$  of the total reaction cross section. Despite this very low production yield investigation of excited states in some of nuclei around  $^{48}\text{Ca}$  was possible due to the high resolving power of the Euroball array.

Scheit et al. [2] observed states in  $^{44,46}\text{Ar}$  using the intermediate-energy Coulomb excitation method and proposed the  $2_1^+$  states of  $^{44}\text{Ar}$  and  $^{46}\text{Ar}$  at  $1144(17) \text{ keV}$  and  $1554(26) \text{ keV}$ , respectively.

We searched for these transitions in the  $\gamma\gamma\gamma$  prompt coincidence cube from  $^{48}\text{Ca} + ^{48}\text{Ca}$  reaction by examining the cross coincidence relationship between complementary Ar and Ti reaction products. Double gates set on prompt transitions in  $^{52}\text{Ti}$  displayed two  $\gamma$ -rays of  $1158.0 \text{ keV}$  and  $1588.4 \text{ keV}$  energy, which were not known previously. The reversed gates showed that both lines are in mutual coincidence and the intensity pattern of prompt yrast transitions from  $^{52}\text{Ti}$  isotope observed in coincidence with the  $1158 \text{ keV}$  and  $1588 \text{ keV}$  lines indicated clearly that these



**Fig. 1.** Coincidence  $\gamma$ -ray spectrum for  $^{44}\text{Ar}$  and  $^{52}\text{Ti}$  complementary products of the  $^{48}\text{Ca}$  (140 MeV) +  $^{48}\text{Ca}$  reaction. It was obtained by summing partial spectra associated with different combinations of double gates on  $^{44}\text{Ar}$  lines: 1158–1588 keV, 1158–693 keV and 1588–693 keV.

$\gamma$ -rays occur in the partner nucleus  $^{44}\text{Ar}$ . In addition, by setting a double gate on the 1158 keV and 1588 keV lines another 692.9 keV  $\gamma$ -transition in  $^{44}\text{Ar}$  was found. A spectrum resulting from the sum of double gates set on the  $^{44}\text{Ar}$  lines is shown in Fig. 1. Examination of intensity relations between 1158, 1588 and 693 keV  $\gamma$ -rays settled clearly the ordering of these transitions, which form a cascade connecting states at 1158, 2746 and 3439 keV excitation energy. There is no doubt that a level at 1158 is the  $2^+$  excitation in  $^{44}\text{Ar}$  observed at 1144(17) keV by Scheit et al. The higher state at 2746 is a candidate for the  $4^+$  excitation. Indeed a level around the same energy was identified at 2.61(16) MeV using the relativistic-energy Coulomb excitation techniques [3] and a  $4^+$  assignment was suggested for it. Finally, the level at 3439 keV is a perfect candidate for the seniority  $2(f_{7/2}^2)6^+$  yrast state.

Similar analysis, performed by employing gates set on known prompt  $\gamma$ -rays in  $^{50}\text{Ti}$ , displayed a 1577 keV  $\gamma$ -ray from a partner nucleus  $^{46}\text{Ar}$ . It is almost certain that this  $\gamma$ -ray is the sought  $2^+ \rightarrow 0^+_{g.s.}$  transition observed by Scheit et al. at 1554(26) keV. This finding settles much

more accurate energy of 1577(1) keV for the  $2^+$  excitation in  $^{46}\text{Ar}$ .

The systematics of known yrast levels in heavy Ar isotopes, including newly placed excitations is shown in Fig. 2.

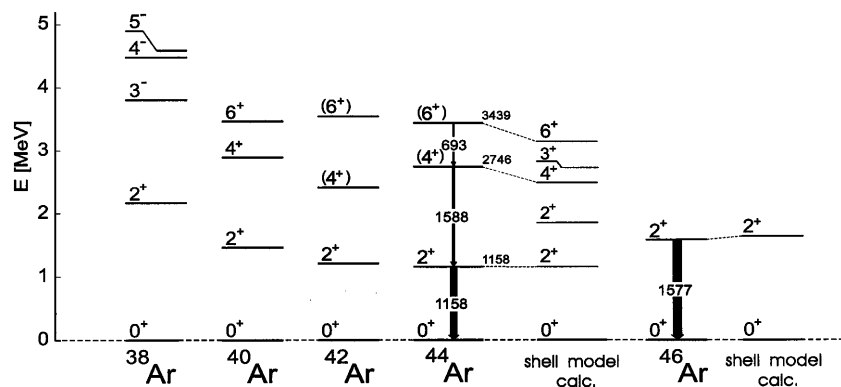
Theoretical description of the properties of neutron-rich *sdfp* shell nuclei was undertaken by Retamosa et al. [6] within the shell model framework. Very good agreement between predictions and experimental results on the  $2^+$  excitation energies and  $B(E2; 0^+_{g.s.} \rightarrow 2^+_1)$  values for  $^{40,42}\text{S}$  and  $^{44,46}\text{Ar}$  was achieved. A conclusion was drawn that, at  $N=28$  in the very neutron-rich regime the magicity persists, even if somewhat eroded by the large neutron excess.

Our new data on higher excitations in  $^{44}\text{Ar}$  offer a more stringent test of these large-scale shell model calculations and interactions used [6]. We have performed such calculations for the  $^{44}\text{Ar}$  and  $^{46}\text{Ar}$  nuclei and their results are compared with the experimental findings in Fig. 2. The agreement is very good for the  $2^+$  states and satisfactory for higher lying yrast levels. In addition, in  $^{44}\text{Ar}$ , the analysis of calculated spectrum and corresponding E2 transitions shows the appearance of a low lying  $\gamma$ -deformed band ( $2^+$ ,  $3^+$ ,  $4^+$  triplet), suggesting gamma softness. Further investigations of this behaviour would be welcome. On the other hand, in  $^{46}\text{Ar}$ , our data confirm the increase of excitation energy of the first  $2^+$  level, which is in line with the persistence of magicity at  $N=28$ . Such phenomenon is indeed predicted by our calculations.

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**Fig. 2.** Systematics of yrast states in neutron-rich Ar isotopes. Arrows indicate the transitions observed in this work. For  $^{44}\text{Ar}$  and  $^{46}\text{Ar}$  the experimental data are compared with the results of the shell model calculations.